Medical Image Analysis

CS 778 / 578

Computer Science and Electrical Engineering Dept. West Virginia University

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Outline

- [The diffusion process](#page-8-0)
- 3 [Perona-Malik : Inhomogeneous diffusion](#page-14-0)
- 4 [Perona-Malik : Anisotropic diffusion](#page-40-0)

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Stability and Convergence

- Stability: noise (from initial conditions, round-off error) is not amplified.
- Convergence: numerical scheme approaches solution of the PDE as $t \rightarrow \infty$

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Convergence of the explicit 1D heat equation

The 1D heat equation, $I_t = I_{xx}$, has solution $I(x, t) = e^{-t} \cos(x)$. This corresponds to the problem with initial condition $I(x, 0) = \cos(x)$.

Discretize only in time (forward)

Observe that $I_{xx}(x, t) = -e^{-t} \cos(x) = -I(x, t)$

$$
\frac{I^{t+\delta}-I^t}{\delta}=-I^t
$$

$$
I^{t+\delta} = I^t - \delta I^t
$$

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Convergence criterion : ratio test

The sequence I^t is convergent if

$$
\lim_{t\to\infty}\left|\frac{I^{t+\delta}}{I^t}\right|<1
$$

The explicit equation we formed earlier

$$
I^{t+\delta} = I^t - \delta I^t
$$

has convergence criterion

$$
\left|\frac{I^{t+\delta}}{I^t}\right| = |1-\delta| < 1
$$

This is satisfied for $0 < \delta < 2$. (Only conditionally convergent.)

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Convergence of the implicit 1D heat equation

Discretize only in time (backward)

$$
\frac{I^{t+\delta}-I^t}{\delta}=-I^{t+\delta}
$$

$$
I^{t+\delta}=I^t-\delta I^{t+\delta}
$$

The implicit equation has convergence criterion

$$
\left|\frac{I^{t+\delta}}{I^t}\right| = \left|\frac{1}{1+\delta}\right| < 1
$$

This is satisfied for $\delta > 0$.

Backward heat equation does not converge in either case.

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Convergence

In general, it can be shown that

- Explicit methods are conditionally convergent.
- Implicit methods are unconditionally convergent.

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Outline

- [The diffusion process](#page-8-0) [Flux](#page-9-0)
	- **[Conservation Laws](#page-11-0)**
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Flux definition

Flux : Rate of movement of *something* per unit area.

What is moving?

- **o** Diffusion: molecules
- Heat: energy

For diffusion the units of flux are $\frac{mol}{m^2s}$

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Flux

Isotropic Diffusion: $j = -d\nabla u$

Fick's First Law : Molecules diffuse from high concentration to low concentration.

- d is a scalar diffusivity constant
- Flux is parallel to concentration gradient, but in opposite direction.

Heat: $q'' = -k\nabla T$ Heat flows from high temperature to low temperature.

Anisotropic Diffusion: $j = -D\nabla u$

Concentration gradient causes a flux which is transformed by the matrix D.

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Discrete example

$$
c_{in} - c_{out} = \Delta c_{stored}
$$

Matter/energy is not created or destroyed.

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Discrete example

$$
(j(x) - j(x + \Delta x))A = \frac{\Delta c_{stored}}{\Delta t}
$$

Matter/energy is not created or destroyed. Extend this example to 2D , and 3D...

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Fick's Second Law

Conservation of Mass

$$
\frac{\partial u}{\partial t} = -\operatorname{div} j
$$

with Fick's First Law $(j = -d\nabla u)$ yields the diffusion equation

$$
\frac{\partial u}{\partial t} = \text{div}(d\nabla u)
$$

- Perona-Malik idea : Make d inhomogeneous $(d(x,y))$
	- ► Slow down / speed up diffusion as needed^{*}
- Weikert idea : Make d anisotropic $(D(x,y))$
	- \triangleright Direct flux as needed

Outline

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[The diffusion process](#page-8-0)

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- [Weaknesses of the standard scale-space paradigm](#page-18-0)
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Scale-space

The need for multiscale image representations: Details in images should only exist over certain ranges of scale.

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Scale-space

Definition: a family of images, $I(x, y, t)$, where

- The scale-space parameter is *t*.
- $I(x, y, 0)$ is the original image.
- Increasing *t* corresponds to coarser resolutions.

 $I(x, y, t)$ can be generated by convolving with wider Gaussian kernels as *t* increases, or equivalently, by solving the heat equation.

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Earlier Scale-space properties

- Causality: coarse details are "caused" by fine details.
- New details should not arise in coarse scale images.
- Smoothing should be homogeneous and isotropic.

This paper will challenge the last property, and propose a more useful scale-space definition.

The new scale-space will be shown to obey the causality property.

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Lost Edge Information

- Edges may disappear.
- Edge location is not preserved across the scale space.
- Region boundaries are blurred.

Gaussian blurring is a local averaging operation. It does not respect natural boundaries.

Linear Scale Space

Def: Scale spaces generated by a linear filtering operation.

- Nonlinear filters, such as the median filter, can be used to generate nonlinear scale-spaces.
- Many nonlinear filters violate one of the scale-space conditions.

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New Criteria

- Causality.
- Immediate localization : edge locations remain fixed.
- Piecewise Smoothing : permit discontinuities at boundaries.

At all scales the image will consist of smooth regions separated by boundaries (edges).

 $\mathbb{R}^n \times \mathbb{R}^n \xrightarrow{\sim} \mathbb{R}^n \times \mathbb{R}^n$

Diffusion equation

$$
\frac{\partial I}{\partial t} = \text{div}(c(x, y, t)\nabla I)
$$

The diffusion coefficient, $c(x, y, t)$ controls the degree of smoothing at each point in *I*.

The basic idea:

Setting $c(x, y, t) = 0$ at region boundaries, and $c(x, y, t) = 1$ at region interior will encourage **intraregion** smoothing, and discourage **interregion** smoothing.

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Conduction coefficient

What properties would we like $c(x, y, t)$ to have?

- \bullet c = 1 at interior of a region.
- \bullet c = 0 at boundary of a region.
- c should be nonnegative everywhere.

Since $c(x, y, t)$ depends on edge information, we need an edge descriptor, $E(x, y, t)$, to compute *c*.

Notation

When written as a function of the edge descriptor, the authors use the symbol g() for conduction coefficient.

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The function $g(||\nabla I||)$

Perona and Malik suggest two possible functions:

$$
g(||\nabla I||) = e^{-\left(\frac{||\nabla I||}{K}\right)^2}
$$

$$
g(||\nabla I||) = \frac{1}{1 + (\frac{||\nabla I||}{K})^{1+\alpha}} \quad (\alpha > 0)
$$

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Effect of varying *K* on $g(||\nabla I||)$

$$
g(||\nabla I||) = \frac{1}{1 + (\frac{||\nabla I||}{K})^{1+\alpha}} \quad (\alpha > 0)
$$

Figure: $K = 2, 4, 6$

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Effect of varying α on $g(||\nabla I||)$

$$
g(||\nabla I||) = \frac{1}{1 + (\frac{||\nabla I||}{K})^{1+\alpha}} \quad (\alpha > 0)
$$

Figure: $\alpha = 1, 3, 5, 7, 9$

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Effect of varying *K* and α on $c(x, y)$

Figure: *I* and $||\nabla I||$.

Effect of varying *K* on $c(x, y)$

Figure: $K = 3, 5, 10, 100$.

As *K* increases, more edges will get smoothed out.

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Effect of varying α on $c()$

Figure: $\alpha = 1, 2, 3, 5$.

As α increases, the cutoff gets sharper.

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Set K every iteration

Compute a histogram, f_i , of $||\nabla I||$

Find *K* such that 90% of the pixels have gradient magnitude $\lt K$. (If $\sum_{i=1}^{b} f_i \geq 0.9n^2$ then bin *b* corresponds to gradient magnitude *K*).

Inhomogeneous diffusion may actually enhance edges, for a certain choice of $c(x, y, t)$.

1D example:

Let
$$
s(x) = \frac{\partial I}{\partial x}
$$
, and $\phi(s) = g(I_x)I_x = g(s)s$.

The 1D inhomogeneous heat equation becomes

$$
I_t = \frac{\partial}{\partial x}(g(I_x)I_x) = \frac{\partial}{\partial x}\phi(s(x))
$$

by chain rule
$$
= \frac{\partial \phi}{\partial s}\frac{\partial s}{\partial x}
$$

$$
I_t = \phi'(s(x))I_{xx}
$$

With a few clever substitutions you can identify the conditions for which ∂ $\frac{\partial}{\partial t}(I_x) > 0$. (See appendix of these notes.)

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Maximum Principle

- The maximum and minimum intensities in the scale-space image $I(x, y, t)$ occur at $t = 0$ (the finest scale image).
- Since new maxima and minima correspond to new image features, the causality requirement of scale-space can satisfied if the evolution equation obeys the maximum principle.
- We will make some less rigorous observations concerning causality...

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Maximum Principle

- Solving the heat equation is equivalent to convolution.
- Convolution is a local averaging operation.
- Averaging is bounded by the values being averaged.

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Maximum Principle

For the Perona-Malik equation

$$
\frac{\partial I}{\partial t} = c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I
$$

Note that at local minima $\nabla I = 0$ and we are evolving by the original heat equation.

It can be shown that this general class of PDEs obeys the maximum principle. See the 2009 class notes for a discussion of the maximum principle for the discretized equations.

Diffusion equation

By the chain rule:

$$
\frac{\partial I}{\partial t} = \text{div}\left(\frac{c(x, y, t)\frac{\partial I}{\partial x}}{c(x, y, t)\frac{\partial I}{\partial y}}\right)
$$
\n
$$
= \frac{\partial c}{\partial x}\frac{\partial I}{\partial x} + c(x, y, t)\frac{\partial^2 I}{\partial x^2} + \frac{\partial c}{\partial y}\frac{\partial I}{\partial y} + c(x, y, t)\frac{\partial^2 I}{\partial y^2}
$$
\n
$$
= c(x, y, t)\nabla^2 I + \nabla c \cdot \nabla I
$$

Notation

The paper uses the symbol Δ to represent the Laplacian. $\Delta I = \nabla^2 I = \text{div}(\nabla I)$

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$$
\frac{\partial I}{\partial t} = c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I
$$

Using centered differences for the Laplacian and gradients:

$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y}(I_{x-1,y}^t + I_{x+1,y}^t + I_{x,y-1}^t + I_{x,y+1}^t - 4I_{x,y}^t) \n+ \frac{(c_{x+1,y} - c_{x-1,y})}{2} (\frac{I_{x+1,y} - I_{x-1,y}}{2}) \n+ \frac{(c_{x,y+1} - c_{x,y-1})}{2} (\frac{I_{x,y+1} - I_{x,y-1}}{2})
$$

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y}(I_{x-1,y}^t + I_{x+1,y}^t + I_{x,y-1}^t + I_{x,y+1}^t - 4I_{x,y}^t) \n+ \frac{(c_{x+1,y} - c_{x-1,y})(I_{x+1,y} - I_{x-1,y})}{2} \n+ \frac{(c_{x,y+1} - c_{x,y-1})}{2}(\frac{I_{x,y+1} - I_{x,y-1}}{2})
$$

• Same diagonal structure as homogeneous heat equation?

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y}(I_{x-1,y}^t + I_{x+1,y}^t + I_{x,y-1}^t + I_{x,y+1}^t - 4I_{x,y}^t) \n+ \frac{(c_{x+1,y} - c_{x-1,y})(I_{x+1,y} - I_{x-1,y})}{2} \n+ \frac{(c_{x,y+1} - c_{x,y-1})}{2}(\frac{I_{x,y+1} - I_{x,y-1}}{2})
$$

- Same diagonal structure as homogeneous heat equation? Yes.
- Symmetric?

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y}(I_{x-1,y}^t + I_{x+1,y}^t + I_{x,y-1}^t + I_{x,y+1}^t - 4I_{x,y}^t) \n+ \frac{(c_{x+1,y} - c_{x-1,y})(I_{x+1,y} - I_{x-1,y})}{2} \n+ \frac{(c_{x,y+1} - c_{x,y-1})}{2}(\frac{I_{x,y+1} - I_{x,y-1}}{2})
$$

- Same diagonal structure as homogeneous heat equation? Yes.
- Symmetric? No.
- Diagonal dominance?

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 $\mathbb{R}^n \times \mathbb{R}^n \xrightarrow{\sim} \mathbb{R}^n \times \mathbb{R}^n$

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y}(I_{x-1,y}^t + I_{x+1,y}^t + I_{x,y-1}^t + I_{x,y+1}^t - 4I_{x,y}^t) \n+ \frac{(c_{x+1,y} - c_{x-1,y})(I_{x+1,y} - I_{x-1,y})}{2} \n+ \frac{(c_{x,y+1} - c_{x,y-1})}{2}(\frac{I_{x,y+1} - I_{x,y-1}}{2})
$$

- Same diagonal structure as homogeneous heat equation? Yes.
- Symmetric? No.
- Diagonal dominance? Data dependent.

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 $\mathbb{R}^n \times \mathbb{R}^n \xrightarrow{\sim} \mathbb{R}^n \times \mathbb{R}^n$

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How do we get from this:

$$
\frac{\partial I}{\partial t} = c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I
$$

to Equation 7?

By splitting the Laplacian and averaging the forward and backward differences in the gradient:

$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = c_{x,y} [(I_{x-1,y}^t - I_{x,y}^t) + (I_{x+1,y}^t - I_{x,y}^t) + (I_{x,y-1}^t - I_{x,y}^t) + I_{x,y-1}^t - I_{x,y}^t)] \n+ \frac{\partial c}{\partial x} [\frac{I_{x+1,y} - I_{x,y}}{2} + \frac{I_{x,y} - I_{x-1,y}}{2}] \n+ \frac{\partial c}{\partial y} [\frac{I_{x,y+1} - I_{x,y}}{2} + \frac{I_{x,y} - I_{x,y-1}}{2}]
$$

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$$
\frac{\partial I}{\partial t} = c(x, y, t) \nabla^2 I + \nabla c \cdot \nabla I
$$

$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = (c_{x,y} - \frac{1}{2} \frac{\partial c}{\partial x}) (I_{x-1,y}^t - I_{x,y}^t) \n+ (c_{x,y} + \frac{1}{2} \frac{\partial c}{\partial x}) (I_{x+1,y}^t - I_{x,y}^t) \n+ (c_{x,y} - \frac{1}{2} \frac{\partial c}{\partial y}) (I_{x,y-1}^t - I_{x,y}^t) \n+ (c_{x,y} + \frac{1}{2} \frac{\partial c}{\partial y}) (I_{x,y+1}^t - I_{x,y}^t)
$$

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These are first order Taylor series approximations

$$
c_{x,y} + \frac{1}{2} \frac{\partial c}{\partial x} \approx c_{x + \frac{1}{2},y}
$$

$$
c_{x,y} - \frac{1}{2} \frac{\partial c}{\partial x} \approx c_{x - \frac{1}{2},y}
$$

$$
c_{x+\frac{1}{2},y} \approx g(\frac{s_{x,y} + s_{x+1,y}}{2})
$$

$$
c_{x-\frac{1}{2},y} \approx g(\frac{s_{x,y} + s_{x-1,y}}{2})
$$

Where $s_{x,y} = ||\nabla I(x, y)||$.

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = g(\frac{s_{x,y} + s_{x-1,y}}{2})(I_{x-1,y}^t - I_{x,y}^t) \n+ g(\frac{s_{x,y} + s_{x+1,y}}{2})(I_{x+1,y}^t - I_{x,y}^t) \n+ g(\frac{s_{x,y} + s_{x,y-1}}{2})(I_{x,y-1}^t - I_{x,y}^t) \n+ g(\frac{s_{x,y} + s_{x,y+1}}{2})(I_{x,y+1}^t - I_{x,y}^t)
$$

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Anisotropic Implementation

Compute $g()$ using the projection of the gradient along one direction. For example, in $g\left(\frac{s_{x,y}+s_{x+1,y}}{2}\right)$ $\frac{2^{3x+1,y}}{2}$), let

$$
s_{x,y} = ||\frac{\partial I}{\partial x}(x,y)||
$$

$$
s_{x+1,y} = ||\frac{\partial I}{\partial x}(x+1,y)||
$$

Computing $s_{x,y}$ using forward differences, and $s_{x+1,y}$ using backward differences

$$
s_{x,y} = ||I_{x+1,y} - I_{x,y}||
$$

$$
s_{x+1,y} = ||I_{x+1,y} - I_{x,y}||,
$$

so $g(\frac{s_{x,y}+s_{x+1,y}}{2})$ $\frac{a_{x+1,y}}{2}$ = $g(||I(x+1,y) - I(x,y)||).$

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$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = g(|I_{x-1,y} - I_{x,y}|)(I_{x-1,y}^t - I_{x,y}^t) \n+ g(|I_{x+1,y} - I_{x,y}|)(I_{x+1,y}^t - I_{x,y}^t) \n+ g(|I_{x,y-1} - I_{x,y}|)(I_{x,y-1}^t - I_{x,y}^t) \n+ g(|I_{x,y+1} - I_{x,y}|)(I_{x,y+1}^t - I_{x,y}^t)
$$

Notation:

The authors use \bigtriangledown to denote finite differences. This is not the gradient operator (∇) .

Book

Image neighborhood system

 $\nabla_N I_{i,j} \equiv I_{i-1,j} - I_{i,j}$ $\nabla sI_{i,j} \equiv I_{i+1,j} - I_{i,j}$ $\bigtriangledown_E I_{i,j} \equiv I_{i,j+1} - I_{i,j}$ $\bigtriangledown_{W}I_{i,j}\ \equiv\ I_{i,j-1}-I_{i,j}$

$$
c_{N_{i,j}} = g(|\nabla_N I_{i,j}|)
$$

\n
$$
c_{S_{i,j}} = g(|\nabla_S I_{i,j}|)
$$

\n
$$
c_{E_{i,j}} = g(|\nabla_E I_{i,j}|)
$$

\n
$$
c_{W_{i,j}} = g(|\nabla_W I_{i,j}|)
$$

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The previous explicit formulation

$$
\frac{I_{x,y}^{t+1} - I_{x,y}^t}{\lambda} = g(|I_{x-1,y} - I_{x,y}|)(I_{x-1,y}^t - I_{x,y}^t) \n+ g(|I_{x+1,y} - I_{x,y}|)(I_{x+1,y}^t - I_{x,y}^t) \n+ g(|I_{x,y-1} - I_{x,y}|)(I_{x,y-1}^t - I_{x,y}^t) \n+ g(|I_{x,y+1} - I_{x,y}|)(I_{x,y+1}^t - I_{x,y}^t)
$$

can be rewritten as

$$
I_{x,y}^{t+1} = I_{x,y}^t + \lambda (c_{N_{i,j}} \nabla_N I_{i,j} + c_{S_{i,j}} \nabla_S I_{i,j} + c_{E_{i,j}} \nabla_E I_{i,j} + c_{W_{i,j}} \nabla_W I_{i,j})^t
$$

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Perona-Malik Implementation

The implementation looks like a discretization of anisotropic diffusion with a diagonal diffusion tensor.

$$
\partial_t u = \text{div}\left(\begin{bmatrix} c_E u_x \\ c_N u_y \end{bmatrix} + \begin{bmatrix} c_W u_x \\ c_S u_y \end{bmatrix}\right)
$$

= $\text{div}\left(\begin{bmatrix} c_E & 0 \\ 0 & c_N \end{bmatrix} \nabla u + \begin{bmatrix} c_W & 0 \\ 0 & c_S \end{bmatrix} \nabla u\right)$

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Discrete Maximum Principle

It can be shown that

- The algorithm will not lead to the production of new local maxima.
- Similarly, no new local minima will be created.
- Therefore, the Perona-Malik algorithm can be used to create scale-space image representations.

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- [Boundary Conditions](#page-52-0)
- [Edge Enhancement](#page-56-0)

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Appendix Boundary Conditions

Constant Boundary Value

$$
(x < 0) \text{ or } (x > n) \rightarrow I(x) = c
$$

For $c = 0$:

$$
I_{xx}(0) \approx -2I(0) + I(1)
$$

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Constant Boundary Slope

Fixing the slope at zero (adiabatic) gives $(x < 0) \to I(x) = I(0)$ $(x > n) \rightarrow I(x) = I(n)$

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 $I_{xx}(0) \approx -I(0) + I(1)$

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Periodic Boundary Conditions

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Reflective Boundary Conditions

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Inhomogeneous diffusion may actually enhance edges, for a certain choice of $c(x, y, t)$.

1D example:

Let
$$
s(x) = \frac{\partial I}{\partial x}
$$
, and $\phi(s) = g(I_x)I_x = g(s)s$.

The 1D inhomogeneous heat equation becomes

I^t =

$$
I_t = \frac{\partial}{\partial x}(g(I_x)I_x) = \frac{\partial}{\partial x}\phi(s(x))
$$

by chain rule
$$
= \frac{\partial \phi}{\partial s}\frac{\partial s}{\partial x}
$$

$$
I_t = \phi'(s(x))I_{xx}
$$

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We are interested in the rate of change of edge slope with respect to time.

$$
\frac{\partial}{\partial t}(I_x) = \frac{\partial}{\partial x}(I_t) \text{ if I is smooth}
$$
\n
$$
= \frac{\partial}{\partial x}(\phi'(s)I_{xx})
$$
\n
$$
= \phi''(s)\frac{\partial s}{\partial x}I_{xx} + \phi'(s)I_{xxx}
$$
\n
$$
= \phi''(s)I_{xx}^2 + \phi'(s)I_{xxx}
$$

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 I, I_x, I_{xx}, I_{xx}

$$
\frac{\partial}{\partial t}(I_x) = \phi''(s(x))I_{xx}^2 + \phi'(s(x))I_{xxx}
$$

For a step edge with $I_x > 0$ look at the inflection point, *p*, where the slope is maximum.

Observe that $I_{xx}(p) = 0$, and $I_{xxx}(p) < 0$.

$$
\frac{\partial}{\partial t}(I_x)(p) = \phi'(s(p))I_{xxx}(p)
$$

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The sign of this quantity depends only on $\phi'(s(p)).$

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At the inflection point:

$$
\frac{\partial}{\partial t}(I_x)(p) = \phi'(s)I_{xxx}(p)
$$

- If $\phi'(s) > 0$, then $\frac{\partial}{\partial t}(I_x)(p) < 0$ (slope is decreasing).
- If $\phi'(s) < 0$, then $\frac{\partial}{\partial t}(I_x)(p) > 0$ (slope is increasing).

Since $\phi(s) = g(s)s$, selecting the function $g(s)$ determines which edges are smoothed and which are sharpened.

 Ω

The function $\phi(s) = g(s)s$

$$
\bullet\ \phi(0)=0
$$

$$
\bullet \ \phi'(s) > 0 \text{ for } s < K
$$

$$
\bullet \ \phi'(s) < 0 \text{ for } s > K
$$

$$
\bullet \ \lim_{s\to\infty}\phi(s)\to 0
$$

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